Growth and Characterization of Graphene/Bismuth Telluride Heterostructures

by Travis M Tumlin, Lily Giri, Mark H Griep, and Shashi P Karna

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Growth and Characterization of Graphene/Bismuth Telluride Heterostructures

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Graphene (G)-based van der Waals heterostructures have recently attracted substantial attention due to the degree of customization afforded with these systems along with the ability to overcome drawbacks associated with their single-layer counterparts. The integration of ultra-high-mobility graphene with other 2-D materials provides an excellent approach for harnessing graphene’s extraordinary electronic properties while allowing for greater complexity and control in switching and energy-harvesting devices. The quality, design, and controlled fabrication of these heterostructures rely on a detailed understanding of the growth mechanisms and transfer techniques. Here we present assembly and characterization of bi- and trilayer G/Bi$_2$Te$_3$/G heterostructures via a combination of transfer and chemical vapor deposition growth techniques. The fabrication of bilayer G/Bi$_2$Te$_3$ was accomplished through epitaxial growth of bismuth telluride nanoplates directly on graphene. The graphene third layer was then deposited using a polymer-assisted transfer technique to create trilayer G/Bi$_2$Te$_3$/G sandwich structures. Electrical measurements of the assembled heterostructure demonstrated an increase in sheet resistance with the growth of Bi$_2$Te$_3$ on graphene. The sheet resistance further decreases with the addition of a top graphene layer below the value of pristine graphene monolayer, suggesting that Bi$_2$Te$_3$ plays an important role in the surface sensitive mobility of electrons across graphene.
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1. Introduction

Graphene has been a focus of intense research in recent years due to its extraordinary electrical, mechanical, and thermal properties.\textsuperscript{1} Ultra-high-electron mobility (up to 200,000 cm\textsuperscript{2}V\textsuperscript{−1}S\textsuperscript{−1}), mechanical strength (1 TPa), thermal conductivity (up to 5.3 × 10\textsuperscript{3} W/mK), and broadband energy absorption make graphene an ideal candidate for next-generation switching and energy harvesting devices.\textsuperscript{2–8} One major shortcoming of pristine graphene is its inherent lack of band gap.\textsuperscript{9,10} This lack of electronic switching (on/off) ability severely limits the electrical applications for graphene in logic devices. Various attempts to engineer a bandgap have focused on the doping of graphene, construction of bilayer graphene, and fabrication of graphene nanoribbons.\textsuperscript{11–13} However, these techniques have their own shortcomings, including degradation of pristine graphene resulting in decreased electron mobility, difficulty in orienting bilayer graphene, and arduous fabrication techniques with repeatability. An alternative to these methods involves the layering of graphene with bandgap-enabled materials such as hexagonal boron nitride (hBN) and transition metal dichalcogenides (TMDs).\textsuperscript{14–18} These techniques are advantageous in that most of these structures can be synthesized through epitaxial growth or through nondestructive layering.\textsuperscript{19–21} A recent advance in graphene-based van der Waals heterostructures is the epitaxial growth of topological insulators (TIs) directly onto graphene.\textsuperscript{22,23} TIs provide an interesting avenue for heterostructure devices due to their unique surface conducting state.\textsuperscript{24,25} Bismuth telluride (Bi\textsubscript{2}Te\textsubscript{3}) is a particularly promising candidate for graphene heterostructures owing to its similar hexagonal symmetry and small bandgap.\textsuperscript{26,27} Photodetector devices based on the epitaxial growth of few-layer bismuth telluride nanoplates on graphene have already been shown to outperform similar devices with other insulating materials.\textsuperscript{28,29} The outstanding performance of these two-layer devices has opened up the possibility for three-layer devices where bismuth telluride is sandwiched between two graphene layers. Devices based on three-layer materials have already been realized for graphene/hBN/graphene and graphene/TMD/graphene systems.\textsuperscript{30–39} To date, however, three-layer graphene/Bi\textsubscript{2}Te\textsubscript{3}/graphene systems have yet to be realized.

In the present work, three-layer graphene/Bi\textsubscript{2}Te\textsubscript{3}/graphene heterostructures were fabricated using a combination of epitaxial growth and polymer-assisted transfer. Bi\textsubscript{2}Te\textsubscript{3} nanoplates show typical hexagonal and triangular shapes with varying thicknesses owing to the van der Waals interactions between the hexagonal unit cell of bismuth telluride and that of graphene. The similarity in the two lattices allows for the epitaxial growth of Bi\textsubscript{2}Te\textsubscript{3} nanoplates on graphene with a seamless interface that promotes effective transfer and separation of carriers. The stacking of
an additional graphene layer is made possible through the use of a polymer-assisted transfer. In-plane sheet resistance shows an increase with the addition of Bi$_2$Te$_3$ over monolayer graphene while the resistance is significantly reduced with the stacking of an additional graphene monolayer. A somewhat insulating nature of bismuth telluride, despite its surface-conducting states, makes it a potential alternative for graphene devices in contrast with other insulating layers such as silicon dioxide (SiO$_2$) or hBN. The present work opens up avenues for future fabrication of graphene-based van der Waals structures for multiple applications.

2. Experimental Procedure

2.1 Graphene Synthesis and Transfer

Graphene was synthesized in a standard chemical vapor deposition (CVD) setup. Prior to deposition, copper foils were electropolished to remove surface impurities and reduce surface roughness. After polishing, copper foils were loaded into a 1-inch quartz tube furnace and evacuated. Hydrogen was then introduced into the system at a pressure of 350 mtorr. Samples were then rapidly heated to 1060 °C and annealed for 30 min in a hydrogen environment. For full-coverage graphene, methane was then introduced into the system at a pressure of 500 mtorr for 20 min. Subsequently, samples were rapidly cooled to room temperature and removed. A polymer-assisted method was used for graphene transfer. Briefly, 1 wt% poly(bisphenol A) carbonate in chloroform was spin coated on a graphene/copper foil sample. Samples were then floated on top of a 0.1-M solution of ammonium persulfate for approximately 2 h until no copper was visible. After which, samples were then rinsed three times in a distilled deionized water bath. Samples were then transferred to the target substrate and allowed to dry in air. Lastly, the adhering polymer was removed by rinsing the sample with chloroform and blow drying with nitrogen.

2.2 Bismuth Telluride Growth

Bismuth telluride was grown directly onto graphene via a vapor transport process. In a typical experiment, silicon or Si/SiO$_2$ substrates with graphene were placed inside a 1-inch quartz tube. A ceramic boat filled with 0.2 mg of bismuth telluride precursor was then placed in the center of the hot zone 9 cm away from the substrates. The tube was then evacuated and flushed three times to remove excess moisture from the system. Then, 500 sccm of argon was introduced with a pressure of 4.43 torr. The furnace was then heated to 460 °C and held at temperature for 20 min. After which, the furnace was shut down and allowed to cool naturally.
2.3 Graphene/Bismuth Telluride Heterostructure Fabrication

Heterostructures were made by taking samples with graphene already transferred onto target substrates and loading into the CVD chamber. A typical growth method was then used to deposit bismuth telluride on graphene. Additional layers of graphene were transferred on top of the bismuth telluride using polymer-assisted transfer.

For nondestructive transfer of the heterostructure sample, the replication method was used. In this method, shown in Fig. 1, a negative replica of the heterostructure film surface was made with cellulose acetate film. Several drops of acetone were placed and spread on the heterostructure sample. Before allowing the acetone to evaporate, the acetate film was placed on one end of the sample and allowed to adhere to the surface via surface tension to prevent the formation of trapped air pockets. Once dry, the acetate film/heterostructure layers were carefully removed from the substrate by pulling from the edge of the film with tweezers. Next, a 4 mm² area of the acetate film/heterostructure negative replica film was cut and carefully placed, heterostructure side down, on transmission electron microscope (TEM) grids. The grids were then placed on a wetted Cu wire mesh in a petri dish. To dissolve the acetate film, acetone was added to the dish and covered with parafilm to prevent evaporation. Acetate film was completely removed after an hour, leaving the heterostructure sample on the TEM grid.

![Fig. 1 Transfer of graphene(G)/bismuth telluride(Bi2Te3)/graphene(G) from silicon growth wafer to TEM grid using cellulose acetate](image)

2.4 Characterization

Optical, laser, and scanning electron microscopy (SEM) were used to determine graphene coverage as well as bismuth telluride domain size and morphology. Raman spectroscopy was used to determine graphene quality throughout the transfer and deposition process. Sheet resistance measurements were performed to see the effect of heterostructure buildup on electronic properties. TEM was used to determine bismuth telluride crystal structure, thickness, and composition. The morphology of the products was analyzed using a Hitachi S4700 field emission scanning electron microscope (FESEM) equipped with an energy-dispersive X-ray system (EDS) used for average compositional analysis. TEM, high-resolution TEM
(HRTEM), and selected area electron diffraction (SAED) of the nanostructures were observed on a JEOL JEM 2100F microscope at 200 kV using an Orius SC1000 camera. Raman spectra for samples were recorded with a Horiba Jobin Yvon LabRam Aramis spectrometer at room temperature using a 532-nm excitation laser wavelength. The scattered light was analyzed with a spectrometer equipped with a detector. Both Raman and EDS results consist of an average of five measurements done at different locations on the samples. Sheet resistance of heterostructures was measured at 10 different locations on the sample using a Signatone S-302-4 in-line probe coupled with a Keithley 6430 SourceMeter.

3. Results

The as-grown G/Bi₂Te₃ heterostructures were imaged and identified by their optical contrast using optical microscopy. Figure 2a shows the optical image of G/Bi₂Te₃ heterostructures on 285-nm-thick SiO₂/Si substrates, with the synthesized nanoplates exhibiting characteristic triangular, hexagonal, and truncated triangular shapes with lateral sizes ranging from a few hundred nanometers up to several microns. The optical contrast between different nanoplates suggests a difference in quintuple layer (QL) thickness. Nanoplates that appear bright white are composed of at least 20 QLs while the less contrasted plates have fewer layers. It is worth noting that individual nanoplates have uniform thickness along the entire plane. Additional synthesis experiments on nucleated graphene domains suggest inhibited growth of Bi₂Te₃ nanoplates on adlayer graphene as well as preferential growth on graphene edges, as shown in Fig. 3. This could be due to defects and dangling bonds present on graphene edges that are not found on the interior.

Figure 2c shows the Raman spectra of the bismuth telluride nanoplates synthesized on graphene. The presence of Bi₂Te₃ is confirmed by the characteristic E₂g and A₁g Raman peaks at 102 cm⁻¹ and 132 cm⁻¹, respectively. The G and 2D peaks along with a lack of the so-called defect peak indicate the presence of pristine graphene in the graphene/Bi₂Te₃ heterostructures. A measured sheet resistance of 158±36 Ω/square, in good agreement with recently reported literature, also verifies the high quality of the CVD synthesized graphene. Figure 2b shows the change in sheet resistance for the heterostructure samples. The G/Bi₂Te₃ sample showed an increase in sheet resistance to 270±44 Ω/square compared to the single-layer graphene while the three-layer G/Bi₂Te₃/G showed a decrease in resistance with a value of 102±6 Ω/square, which is lower than that of single-layer graphene. The possibilities for this phenomena are discussed later in the report.
Fig. 2  Characterization of G/Bi$_2$Te$_3$/G heterostructures: a) Difference in optical contrast for each structure on SiO$_2$/Si. b) Sheet resistance of monolayer graphene compared to G/Bi$_2$Te$_3$ and G/Bi$_2$Te$_3$/G heterostructures. The connecting lines between data points are intended as a guide for the eye. c) Raman spectrum for various stacking orders.
The crystal structure and chemical composition of the as-synthesized Bi$_2$Te$_3$ was further studied using TEM. The samples were mechanically transferred from the Si wafer to the TEM grid using 0.1-M polymethyl methacrylate in acetone. Figure 4a shows a regular TEM image of a single triangular nanoplate. High-resolution TEM in Fig. 4b shows expected hexagonal lattice fringes with a lattice spacing of 0.22 nm corresponding to (110) plane. The corresponding SAED pattern exhibits a sixfold symmetry, which confirms the single crystalline nature of the Bi$_2$Te$_3$ nanoplate. The d spacing (0.22 nm) calculated from the diffraction spots conclusively proves that those are due to diffraction from (110) plane of Bi$_2$Te$_3$ (R3m) phase. The thickness of a single hexagonal nanoplate was calculated by electron energy loss spectroscopy (EELS) and shown in Fig. 4c. The thickness profile calculated along the rectangular section in Fig. 4c was approximately 15 nm, which corresponds to approximately 15 QLs. The EDS spectrum shown in Fig. 4d gives information about the elements present as well as the composition of the nanoplates. The presence of only Bi and Te suggests that the vapor transport–deposited nanoplates are highly pure. The Cu and the C lines are coming from the
TEM grid and the holey carbon support film. Compositional analysis confirms the composition of the nanoplates with a Bi:Te atomic ratio of 2:3, indicating the Bi$_2$Te$_3$ phase.

Fig. 4  a) TEM image of a Bi$_2$Te$_3$ triangular nanoplate, b) HRTEM of Bi$_2$Te$_3$ showing lattice fringes and SAED pattern, c) thickness calculation by EELS, and d) EDS spectrum

Three-layer heterostructures were fabricated using well-established polymer-assisted graphene transfer techniques.$^{48-50}$ Figure 5a shows the contrast difference between two types of nanostructures on SiO$_2$/Si: G/Bi$_2$Te$_3$, and G/Bi$_2$Te$_3$/G. SEM in Fig. 5b shows the two-layer G/Bi$_2$Te$_3$ compared with the three-layer G/Bi$_2$Te$_3$/G heterostructure. Folds and wrinkles in the images suggest the presence of a thin graphene membrane covering the Bi$_2$Te$_3$ nanoplates. The SEM images of Bi$_2$Te$_3$ nanoplates on graphene and graphene/Bi$_2$Te$_3$/graphene are shown in Fig. 5c and d, respectively. Figure 5c shows a clear contrast between the bare Si and the graphene with regular-shaped Bi$_2$Te$_3$ nanoplates. The SEM image of the graphene/Bi$_2$Te$_3$/graphene heterostructure (Fig. 5d) clearly shows folds and wrinkles of a second graphene layer on top of the Bi$_2$Te$_3$ nanoplates. This is commonly observed in large-area CVD growth of graphene on metals due to the difference in the thermal expansion coefficients of the two dissimilar systems.$^{51-54}$ The presence of wrinkles confirms the formation of the G/Bi$_2$Te$_3$/G heterostructure.
Fig. 5 Characterization of G/Bi₂Te₃ heterostructures: a) Optical micrograph of G/Bi₂Te₃ transferred onto 285-nm SiO₂/Si. b) SEM of G/Bi₂Te₃ and G/Bi₂Te₃/G heterostructures separated by a thick dark line running through the middle of the image. c) Magnified image of G/Bi₂Te₃ bilayer. d) SEM of G/Bi₂Te₃/G trilayer heterostructure with prominent wrinkles (bright lines) in the top layer graphene.

The microstructure and the composition of the heterostructures were further investigated using TEM. To prevent wrinkling, a replication transfer method was used for TEM sample preparation. Figure 6a shows the regular TEM image of the Bi₂Te₃ nanoplates grown on graphene. Randomly oriented hexagonal and triangular shaped particles can be clearly seen on the graphene film supported by the holey carbon film of the TEM grid. The particles, with lateral sizes varying from a couple hundred nanometers to several microns, are observed only on the graphene film, implying the preferential growth of Bi₂Te₃ on graphene. The SAED pattern was taken from a random area of the sample and shown in Fig. 6b. The pattern reveals hexagonal symmetry from the Bi₂Te₃ as well as that of graphene. The d spacings were calculated by rotational average method and agree very well with graphene (0.35 nm) and (110) plane spacing of Bi₂Te₃ (0.22 nm), implying the presence of pristine-quality graphene and highly pure single crystalline Bi₂Te₃.
For higher-contrast imaging, scanning transmission electron microscopy (STEM) was used to further observe the Bi$_2$Te$_3$ nanoplates on graphene. In the dark field image obtained by this method, the intensity varies with both the atomic number and the thickness of the material so that higher atomic number and thicker material will appear brighter. As observed from the Fig. 6c, the contrast between the bright Bi$_2$Te$_3$ nanoplates grown on graphene film and the holey carbon films is very prominent and more clearly visible. The STEM image of trilayer G/Bi$_2$Te$_3$/G (Fig. 6d) clearly shows folds and the wrinkles, which arise from the presence of a top graphene layer on Bi$_2$Te$_3$ nanoplates.

The sheet resistance, measured at room temperature, of the monolayer graphene and the G/Bi$_2$Te$_3$ and G/Bi$_2$Te$_3$/G heterostructures, shown in Fig. 2b, exhibits an interesting trend, which can be used as a useful guide to further tailor the electronic properties of G/Bi$_3$Te$_3$ heterostructures. The sheet resistance of monolayer graphene is measured to be approximately 150 $\Omega$/square, which compares well with the previously reported values.\textsuperscript{55-57} The sheet resistance of the G/Bi$_2$Te$_3$ is much higher than the monolayer graphene, which is understandable, since the latter has a much lower conductivity than graphene. However, the addition of another graphene monolayer on top of the bilayer G/Bi$_2$Te$_3$ considerably reduces the sheet resistance of the top G layer. In fact, the measured value for the top G layer in the trilayer
G/Bi₂Te₃/G heterostructure is even lower than the pristine graphene monolayer. This considerable reduction in sheet resistance of G in the trilayer G/Bi₂Te₃/G could potentially result from (a) reduced surface defects/scattering centers on the second graphene layer compared to the single layer graphene,⁵⁸ (b) enhanced density of states due to coupling between the Bi₂Te₃ surface states and graphene states at the Fermi level, and/or (c) inter-layer coupling between the states of the two graphene monolayers. Enhanced density of states in bilayer graphene compared to a single monolayer, observed experimentally and predicted by density functional theory calculations, have been shown to lead to reduced resistivity.⁵⁹,⁶⁰ This could very well be the case in the present study, since Bi₂Te₃ nanostructures are not uniformly deposited on the bottom graphene surface, allowing graphene-graphene coupling in the regions lacking Bi₂Te₃ structures. Thus, mechanism (a) and (c) appear to be more plausible, although we cannot rule out the coupling between Bi₂Te₃ surface states and the top graphene states. Further experiments are underway to understand the underlying mechanism, including the role of trace amounts of polymer residues and other impurities and humidity on electronic properties of G/Bi₂Te₃. Van der Waals heterostructures will be reported in subsequent communications.

4. Conclusion

In this work, graphene/Bi₂Te₃ heterostructures were fabricated through a vapor deposition and polymer-assisted transfer process. Epitaxial growth of Bi₂Te₃ on graphene showed well-ordered nanoplates with lateral dimensions of up to approximately 10 µm and thicknesses on the order of several to tens of nanometers. Diffraction and Raman studies confirmed the high quality and crystallinity of the stacked heterostructures. Electrical characterizations show an increase and subsequent decrease in sheet resistance in going from a pristine graphene monolayer to bilayer G/Bi₂Te₃ to trilayer G/Bi₂Te₃/G heterostructures, suggesting a prominent role of Bi₂Te₃ in modulating electronic property of G-based van der Waals solids. Although the actual mechanism for the observed trend in the electrical property of the heterostructures is not clearly understood, this work provides a basis for further tailoring the electronic properties of G/Bi₂Te₃ heterostructures for useful applications.
5. References


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<td>Bi$_2$Te$_3$</td>
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<td>CVD</td>
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